

Additional Homeworks. Gravitation and Cosmology. Year 2017/2018

1. Consider the following space-time

$$ds^2 = \frac{a^2}{x^2}(-dt^2 + dx^2), \quad t \in \mathbb{R}, \quad x \in \mathbb{R}^+.$$

Compute all the geodesics (space, null and time-like ones).

2. Compute all the geodesics of the following metric

$$ds^2 = (dx^1)^2 - \frac{1}{(x^2)^2}(dx^2)^2.$$

3. Show

$$V_{\alpha;\nu\kappa} - V_{\alpha;\kappa\nu} = V_{\sigma}R^{\sigma}{}_{\alpha\nu\kappa}.$$

4. A Killing vector ξ satisfies $\mathcal{L}_{\xi}\mathbf{g} = 0$, where \mathcal{L}_{ξ} is the Lie-derivative along the vector ξ and \mathbf{g} is the metric. Show that

$$\xi_{(\alpha;\beta)} = 0.$$

Build some metrics with explicit symmetries and check the previous equation.

5. Let ξ_1 and ξ_2 be two Killing vectors. Show that

- $[\xi_1, \xi_2]$ is a Killing vector.
- $\lambda_1\xi_1 + \lambda_2\xi_2$ is also a Killing vector ($\lambda_1, \lambda_2 \in \mathbb{R}$).
- Check the results (a) and (b) by building some metrics with symmetries.

6. (a) Show that the interval ds^2 of a plane space-time (t, x, y, z) which is rotating with angular speed Ω about the z axis of an inertial frame can be written as

$$ds^2 = -[1 - \Omega^2(x^2 + y^2)]dt^2 + 2\Omega(ydx - xdy)dt + dx^2 + dy^2 + dz^2.$$

- Compute the differential equations which follow the time-like geodesics.
- Show that the Newtonian limit of the time-like geodesics are the Newton equations written in a non-inertial frame. Identify and name all the terms.

7. Consider a planet orbiting in the equatorial plane of the Sun.

- Assuming the Sun is oblate, compute its gravitational potential to order $O(1/r^3)$ in the equatorial plane.
- Write the perturbed Kepler problem and solve it perturbatively to show that the perihelion precesses.
- Consider Mercury as an example and compute for it the precession angular speed of the perihelion.
- Find in the literature the Sun's oblateness and compare your result with the measured Mercury precession.

Hint: Check, for example, Goldstein's book.

8. Consider a planet orbiting the Sun in presence of interplanetary dust. One can show that the total potential (Sun+dust) is (assuming a spherical and uniform distribution of the dust):

$$V(r) = -\frac{GM}{r} + \frac{1}{2}Cr^2,$$

where M is the solar mass and C is a positive constant.

- (a) Could you relate the constant C to the density of interplanetary dust?
 (b) Write the perturbed Kepler problem and solve it perturbatively ($C \ll GM/l^3$, l is a measure of the size of the Mercury orbit) to show that the perihelion precesses.
 (c) Compute the precession angular speed of the perihelion.

Hint: Check, for example, Goldstein's book.

9. Consider the following metric:

$$ds^2 = dudv + \log(x^2 + y^2)du^2 - dx^2 - dy^2$$

with $0 < x^2 + y^2 < 1$. Write the equations of the geodesics and show that $K = x\dot{y} - y\dot{x}$ is a conserved quantity.

10. (a) Let χ^μ be a vector satisfying

$$\nabla_\mu \chi_\nu + \nabla_\nu \chi_\mu = cg_{\mu\nu}$$

being c a positive constant and let γ be a null geodesic with tangent vector u^μ . Show that $u^\mu \chi_\mu$ is constant along the geodesic γ .

- (b) Let $T_{\mu\nu}$ and ξ be a symmetric conserved tensor ($\nabla_\mu T^{\mu\nu} = 0$) and a Killing vector respectively. Show that the current defined as $J^\mu \equiv T^{\mu\nu} \xi_\nu$ is conserved ($\nabla_\mu J^\mu = 0$).

Hint: A Killing vector satisfies $\xi_{(\alpha;\beta)} = 0$ (see Problem 4).

11. Compute $a^\mu = Du^\mu/D\tau$ for a stationary particle in the Schwarzschild geometry. Compute its norm (a^2). Show that the force needed to hold the particle at rest is infinite as $r = 2M$.

12. Consider the following 4 + 1 metric

$$ds_5^2 = e^{2A(v)} \eta_{\alpha\beta} dx^\alpha dx^\beta + dv^2$$

where v is the fourth spatial coordinate. Find the differential equation which is satisfied by $A(v)$ if the metric is solution of the Einstein equations in the vacuum with cosmological constant ($\Lambda \neq 0$). What is the sign of Λ ? What happens if $\Lambda = 0$?

13. Consider a 4 + 1 space-time.

- What is the Newtonian gravitational potential?
- Show that the metric

$$ds^2 = - \left(1 - \frac{r_s^2}{r^2} \right) dt^2 + \left(1 - \frac{r_s^2}{r^2} \right)^{-1} dr^2 + r^2 d\Omega_3^2$$

is a solution of the Einstein equations in the vacuum (is the 4 + 1 analogous of the Schwarzschild solution). Moreover, $d\Omega_3^2$ is the metric of the 3-sphere S^3 and r_s is a constant.

14. Find the function $f(v)$ so that the metric

$$ds^2 = 2dudv + f(v)dw^2 + 2dwdz$$

is solution of the Einstein equations in the vacuum. Determine the conditions $f(v)$ should satisfy in order to have a plane space-time.

15. Compute the non zero components (modulo symmetries) of the Riemann tensor in the Schwarzschild geometry. Write these Riemann tensor components in the associated orthonormal basis.

Consider a radially falling observer in a Schwarzschild black hole. Write the geodesic deviation equations in the orthonormal basis associated to the falling observer.

Hint: Perform an (instantaneous) Lorentz transformation between the orthonormal basis of an observer at infinite and that of the falling observer.

16. Consider the following metric

$$ds^2 = -(1 - G(x, y))dt^2 - 2G(x, y)tdtz + (1 + G(x, y))dz^2 + dx^2 + dy^2.$$

- Write the geodesic equations and the Christoffel's symbols.
- Identify two Killing vectors and write the two associated conserved quantities. Compute the norms of these Killing vectors (time, null or spacelike?).
- Assuming vacuum, write the Einstein equations to obtain differential equations for the function G and find some solution for G .
- If $G(x, y) = \Phi(x^2 + y^2)$, find a third Killing vector, and write an explicit solution for the function G .

17. Consider two observers at rest in a gravitational field given by the following metric (anti-de Sitter space):

$$ds^2 = -\left(1 + \frac{r^2}{l^2}\right)dt^2 + \frac{dr^2}{1 + \frac{r^2}{l^2}} + r^2d\Omega_2^2,$$

l being a parameter with dimensions of length. The first observer has spatial coordinates $r_1 = a$, $\theta_1 = \pi/2$, $\phi_1 = 0$ and the second $r_2 = b$, $\theta_2 = \pi/2$, $\phi_2 = 0$, with $b \ll a \ll l$. The first observer emits light which propagates in the direction $\theta = \pi/2$ and $\phi = 0$. This light is detected by the second observer. What is the redshift as seen by the second observer (blue or red)?

18. The metric of Reissner-Nordström

$$ds^2 = -\left(1 - \frac{2M}{r} + \frac{Q^2}{r^2}\right)dt^2 + \left(1 - \frac{2M}{r} + \frac{Q^2}{r^2}\right)^{-1}dr^2 + r^2d\Omega^2,$$

is a solution of the Einstein equations with $\Lambda = 0$ and the energy-momentum tensor produced by an electromagnetic field. This metric describes the external gravitational field of a charged star or a static charged black hole. The parameters M and Q are related with the mass and charge of the star or black hole. We will assume hereafter that $M^2 > Q^2$, why?

- Show that a chargeless massive particle falling radially can never reach $r = 0$. Compare this situation with a particle falling in the Schwarzschild metric.
- Consider a particle initially is at rest at infinity. Compute the minimum radial coordinate of the particle during its falling, r_{\min} . Check that $r_{\min} < r_-$, where r_- is the smallest root of the equation

$$1 - \frac{2M}{r} + \frac{Q^2}{r^2} = 0.$$

19. Consider a Schwarzschild black hole with mass M .

- Show that the light can follow a circular orbit with radial coordinate $r = 3M$.
- Discuss the stability of this orbit.
- Compute in coordinate time the period of this orbit.
- Compute the period measured by an observer at rest in the infinity.
- Consider an observer at rest in a fixed point of the orbit $r = 3M$: compute the period measured by this observer?